An Augmented Plane Wave Plus Local Orbital Program for Calculating Crystal Properties

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http://www.wien2k.at
WIEN2k consists of many independent F90 programs, which are linked together via C-shell scripts.

Each „case“ runs in his own directory ./.case

The „master input“ is called case.struct

Initialize a calculation: init_lapw

Run scf-cycle: run_lapw (runsp_lapw)

You can run WIEN2k using any www-browser and the w2web interface, but also at the command line in an xterm.

Input/output/scf files have endings as the corresponding programs:

case.output1..lapw1; case.in2..lapw2; case.scf0..lapw0

Inputs are generated using STRUCTGEN(w2web) and init_lapw
**w2web: the web-based GUI of WIEN2k**

- **Based on www**
  - **WIEN2k can be managed remotely via w2web**

- **Important steps:**
  - **start w2web on all your hosts**
    - login to the desired host (ssh)
    - w2web (at first startup you will be asked for username/password, port-number, (master-)hostname. creates ~/.w2web directory)
  - **use your browser and connect to the (master) host:portnumber**
    - firefox http://fp98.zserv:10000
  - **create a new session on the desired host** (or select an old one)
w2web GUI (graphical user interface)

- **Structure generator**
  - spacegroup selection
  - import cif or xyz file
  - step by step initialization
    - symmetry detection
    - automatic input generation

- **SCF calculations**
  - Magnetism (spin-polarization)
  - Spin-orbit coupling
  - Forces (automatic geometry optimization)

- **Guided Tasks**
  - Energy band structure
  - DOS
  - Electron density
  - X-ray spectra
  - Optics
Spacegroup P4$_2$/mm

Structure given by:

- spacegroup
- lattice parameter
- positions of atoms (basis)

Rutile TiO$_2$:
P4$_2$/mm (136)
a = 8.68, c = 5.59 bohr
Ti: (0,0,0)
O: (0.304,0.304,0)
Structure generator

- Specify:
  - Number of nonequivalent atoms
  - lattice type (P, F, B, H, CXY, CXZ, CYZ) or spacegroup symbol
    - if existing, you must use a SG-setting with inversion symmetry:
      - Si: $\pm(1/8,1/8,1/8)$, not $(0,0,0)+(1/4,1/4,1/4)$!
  - lattice parameters $a, b, c$ (in Å or bohr)
  - name of atoms (Si) and fractional coordinates (position)
    - as numbers (0.123); fractions (1/3); simple expressions (x-1/2,...)
    - in fcc (bcc) specify just one atom, not the others in (1/2,1/2,0; ...)

- „save structure“
  - updates automatically $Z$, $r_0$, equivalent positions

- „set RMT and continue“: (specify proper “reduction” of NN-distances)
  - non-overlapping „as large as possible“ (saves time), but not larger than 2.5 bohr
  - RMT for sp (d) - elements 10-20 % smaller than for d (f) elements
  - largest spheres not more than 50 % larger than smallest sphere
  - Exception: H in C-H or O-H bonds: $RMT \sim 0.6$ bohr ($RKMAX \sim 3-4$)
  - Do not change RMT in a „series“ of calculations, RMT equal for same atoms

- „save structure – save+cleanup“
Program structure of WIEN2k

- **init_lapw**
  - step-by-step or batch initialization
  - symmetry detection (F, I, C-centering, inversion)
  - input generation with recommended defaults
  - quality (and computing time) depends on **k-mesh** and **R.Kmax** (determines #PW)

- **run_lapw**
  - scf-cycle
  - optional with SO and/or LDA+U
  - different convergence criteria (energy, charge, forces)

- **save_lapw tic_gga_100k_rk7_vol0**
  - cp case.struct and clmsum files,
  - mv case.scf file
  - rm case.broyd* files
The convergence criterion in APW is the product of $R_{MT} \cdot \text{Kmax}$

$$\Psi = \sum_{K_n}^{K_{MAX}} c_{K_n} e^{iK_n r}$$

http://www.wien2k.at/reg_user/faq/rkmax.html

medium quality convergence for **smallest** atom:

- basis set scales with $\text{RKmax}^3$
- cputime scales with $\text{N}_{PW}^3$

increasing Rkmax by 10 % doubles cputime

START with **SMALL** Rkmax (relaxation), *increase/test* later
BZ integration, “FERMI”-methods

- Replace the “integral” of the BZ by a finite summation on a mesh of “k-points”
  \[ \rho(r) = \sum_{n} \int_{E_n < E_F} \psi_{k,n}^* \psi_{k,n} \, d^3k = \sum_{k,n} w_{k,n} \psi_{k,n}^* \psi_{k,n} \]

- weights \( w_{k,n} \) depend on \( k \) and bandindex \( n \) (occupation)
  - for full “bands” the weight is given by “symmetry”
    - \( w(\Gamma) = 1, \ w(x) = 2, \ w(\Delta) = 4, \ w(k) = 8 \)
  - shifted “Monkhorst-Pack” mesh

- for partially filled bands (metals) one must find the Fermi-energy (integration up to NE) and determine the weights for each state \( E_{k,n} \)
  - linear tetrahedron method (TETRA, eval=999)
  - linear tetrahedron method + “Bloechl” corrections (TETRA)
  - “broadening methods”
    - gauss-broadening (GAUSS 0.005)
    - temperature broadening (TEMP/TEMPS 0.005)

- broadening useful to damp scf oscillations, but dangerous (magnetic moment)
**k-mesh generation**

- **x kgen** (generates k-mesh and reduces to irreducible wedge using symmetry)
  - automatically “adds inversion”
    - time inversion holds and $E(k) = E(-k)$
    - except in magnetic spin-orbit calculations (x -so kgen; uses case.ksym file)
  - x -fbz kgen (generates “full mesh” in BZ)

- **always “shift” the mesh for scf-cycle**
  - gaps often at $\Gamma$ ! (might not be in your mesh)

- **small unit cells and metals** require large k-mesh (1000-100000)
- **large unit cells and insulators** need only 1-10 k-points
- **use at first a fairly coarse mesh** for scf/relaxations
- **continue later with finer mesh**
  - mesh was good if nothing changes and scf terminates after few (3) iterations
- **use even finer meshes** for DOS, spectra, optics,…
Program execution:

- All programs are executed via the „master“ shell-script „x“:
  
  ```
  x lapw2 -up -c
  ```

- This generates a „def“ file: `lapw2.def`
  ```
  5,'tin.in2c', 'old', 'formatted'
  6,'tin.output2up', 'unknown','formatted'
  8,'tin.clmvalup', 'unknown','formatted'
  10,'./tin.vectorup','unknown','unformatted'
  ```

- and executes: `lapw2c lapw2.def`

- All WIEN2k-shell scripts have long and short names:
  ```
  x_lapw; runsp_lapw, runfsm_lapw ➔ x; runsp; runfsm
  ```

- All scripts have a „help“ switch „-h“, which explains flags and options (without actually execution)
  ```
  x -h          x lapw1 -h
  ```
run_lapw [options] (for nonmagnetic cases)

- ec 0.0001 convergence of total energy (Ry)
- cc 0.0001 convergence of charge distance (\text{e}^{-})
- fc 1.0 convergence of forces (mRy/bohr)
- it (-it1,-it2, -noHinv) iterative diagonalization (large speedup)
- p parallel calculation (needs .machines file)
- so add spin-orbit (only after „init_so‟)

Spacegroups without inversion use automatically lapw1c, lapw2c (case.in1c,in2c)

case.scf: master output file, contains history of the scf-cycle

most information is stored with some „labels‟ (grep :label case.scf)

<table>
<thead>
<tr>
<th>:ENE</th>
<th>:DIS</th>
<th>:FER</th>
<th>:GAP</th>
<th>:CTO001</th>
<th>:NTO001</th>
<th>:QTL001</th>
</tr>
</thead>
<tbody>
<tr>
<td>:FOR002: 2.ATOM</td>
<td>19.470</td>
<td>0.000</td>
<td>0.000</td>
<td>19.470</td>
<td></td>
<td></td>
</tr>
<tr>
<td>:FGL002: 2.ATOM</td>
<td>13.767</td>
<td>13.767</td>
<td>0.000</td>
<td>total forces</td>
<td></td>
<td></td>
</tr>
<tr>
<td>:LAT</td>
<td>:VOL</td>
<td>:POSxxx</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
</tbody>
</table>
Getting help

- `_lapw -h` „help switch“ of all WIEN2k-scripts
- `help_lapw`:
  - opens `usersguide.pdf`; Use `^f keyword to search for an item („index“)`
- html-version of the UG: (`$WIENROOT/SRC_usersguide/usersguide.html`)
- `http://www.wien2k.at/reg_user`
  - FAQ page with answers to common questions
  - Update information: When you think the program has an error, please check newest version
  - Textbook section: *DFT and the family of LAPW methods by S. Cottenier*
- Mailing-list:
  - subscribe to the list (always use the same email)
  - full text search of the „digest“ (your questions may have been answered before)
  - posting questions: Provide sufficient information, locate your problem (case.dayfile, *.error, case.scf, case.outputX).
  - „My calculation crashed. Please help.“ This will most likely not be answered.
most common problems

- "QTL-B" value too large - STOP (or :WARN): “ghostbands”
  - identify for which eigenvalue, atom and ℓ it happens, check $E_F$
    (case.scf2, case.output2)
  - identify the corresponding linearization energies in case.scf1
  - change the corresponding linearization energy in case.in1
    - compare and check with :EPL and :EPH lines in case.scf2
    - default E-parameters are adapted automatically but may need changes for
      - surfaces, molecules (negative EF) or heavy elements (EF often larger than 1.0)
      - add a local orbital (or adjust its energy)
  - if QTL-B occurs for an atom with large RMT, reduce RMT
    - this may happen for larger RKMAX ("numerical linear dependency")

- scf-cycle diverges (grep :DIS case.scf):
  - check structure (most likely a wrong structure caused divergence);
  - reduce mixing in case.inm slightly; rm *.broyd* case.scf; x dstart
  - check E-parameters (see above), check :NEC01
case.in1

- WFFIL: EF=0.634 (WFPR, SUPWF)
- 7.00 10 4 (R-MT*K-MAX; MAX L IN WF, V-NMT)
- 0.30 5 0, global E-param with N other, napw
- 0 0.30 0.000 CONT 1 Es
- 0 -3.72 0.005 STOP 1 Es-LO with search
- 1 -2.07 0.010 CONT 1 Ep with search
- 1 0.30 0.000 CONT 1 Ep-LO
- 2 0.30 0.010 CONT 1 0/1..LAPW/APW+lo

- K-VECTORS FROM UNIT: 4

\[ \Psi = \sum_{n}^{KMAX} c_{Kn} e^{iK_{n}r} \]

\[ \Phi_{Kn} = \sum_{l}^{l_{max}} A_{lm} u_{l}(E_{l}, r)Y_{lm} \]

\[ H_{NS}^{LM} = \langle \Phi_{l}^{LM} | V_{NS}^{LM} | \Phi_{l}^{LM} \rangle \]
case.klist, case.in2

- **GAMMA**
  - 0 0 0 40 1.0 IX, IY, IZ, IDIV, WEIGHT
  - 1 0 0 40 6.0
- ...
- **X**
  - 40 0 0 40 3.0
- **END**

**case.in2:**

- **TOT**
  - (TOT, FOR, QTL, EFG, FERMI)
  - -9.0 16.0 0.50 0.05 EMIN, NE, ESEPARMIN, ESEPAR0
- **TETRA**
  - 0.000 (GAUSS, ROOT, TEMP, TETRA, ALL eval)
  - 0 0 4 0 4 4 6 0 6 4
  - 0 0 4 0 4 4 6 0 6 4
- **FILE**
  - FILE/Nofile write recprlist

\[
\rho(r) = \sum_{LM} \rho_{LM}(r) Y_{LM}(\hat{r})
\]

\[
\rho(r) = \sum_{G} \rho_{G} e^{iGr}
\]
Properties with WIEN2k - I

- **Energy bands**
  - *classification of irreducible representations*
  - ´character-plot´ (emphasize a certain band-character)

- **Density of states**
  - *including partial DOS with l and m- character (eg. p_x, p_y, p_z)*

- **Electron density, potential**
  - *total-, valence-, difference-, spin-densities, \( \rho \) of selected states*
  - 1-D, 2D- and 3D-plots (Xcrysden)
  - *X-ray structure factors*
  - *Bader’s atom-in-molecule analysis, critical-points, atomic basins and charges* (\( \nabla \rho \cdot \vec{n} = 0 \))
  - *spin+orbital magnetic moments (spin-orbit / LDA+U)*

- **Hyperfine parameters**
  - *hyperfine fields (contact + dipolar + orbital contribution)*
  - *Isomer shift*
  - *Electric field gradients*
partial charges “qt!l” + DOS

- be sure to have case.vector on a dense tetrahedral mesh after a scf calculation
  - eventually:
    - x kgen
    - edit case.in1 (larger Emax)
    - x lapw1

- case.outputt
  - integrated DOS

- case.dos1ev (3ev)
  - text-file for plotting
  - $E$-zero at $E_F$

Density of states

Session: TiC
/susi/pbaha/lapw/TiC

**Density of states**

- x lapw2 -qt Calculate partial charges ☑ interactively
- edit TiC.int Edit input-file for TETRA
- x tetra Calculate partial DOS ☑ interactively
- edit TiC.output Check output of TETRA
- dosplot Plot DOS

**File:**
/susi/pbaha/lapw/TiC/TiC.int

**Header from TiC.qtl:**

| ATOM | 1  | tot.0.1.2.3.xdos(i,j);i=1-3;j=1-3|xldos2) |
|------|----|----------------------------------|
| ATOM | 2  | tot.0.1.2.3.xeg;j=1-3            |

**Title**

<table>
<thead>
<tr>
<th>3</th>
<th>0.50</th>
<th>0.002</th>
<th>1.500</th>
<th>0.003</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>total</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>0</td>
<td>1</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>2</td>
<td>2</td>
<td>Atom1-s</td>
<td></td>
<td></td>
</tr>
<tr>
<td>2</td>
<td>3</td>
<td>Atom2-s</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>
partial charges:

- **local rotation matrix:**
  - transfers z (y) into highest symmetry
  - reduces terms in LM series
  - “chemical” interpretation
    - p\(_x\) is different from p\(_y\)

- see case.struct and case.outputs

- **x qtl** (instead of \(x\) lapw2 -qtl)
  - **f-orbitals**
  - **qtls for different coordinate system** (eg. “octahedral” in TiO\(_2\))
  - relativistic basis (\(p_{1/2}-p_{3/2}\) or \(d_{3/2}-d_{5/2}\) splitting in so calculation)
  - for angular dependend TELNES (ISPLIT 88, 99)
Properties with WIEN2k - I

- Energy bands
  - classification of irreducible representations
  - ´character-plot´ (emphasize a certain band-character)
- Density of states
  - including partial DOS with l and m- character (eg. \( p_x \), \( p_y \), \( p_z \))
- Electron density, potential
  - total-, valence-, difference-, spin-densities, \( \rho \) of selected states
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  - X-ray structure factors
    - Bader’s atom-in-molecule analysis, critical-points, atomic basins and charges \( (\nabla \rho \vec{n} = 0) \)
    - spin+orbital magnetic moments (spin-orbit / LDA+U)
- Hyperfine parameters
  - hyperfine fields (contact + dipolar + orbital contribution)
  - Isomer shift
  - Electric field gradients
  - NMR chemical shifts
Atoms in Molecules

- Theory to characterize atoms and chemical bonds from the topology of the electron density, by R.F. Bader
  (http://www.chemistry.mcmaster.ca/faculty/bader/aim/aim_0.html)

Electron density of $C_2H_4$
Bonds are characterized by „critical points“, where $\nabla \rho = 0$

- density maximum: (3,-3); 3 negative curvatures $\lambda$, (at nucleus or non-NM)
- bond CP: (3,-1): 2 negative, 1 positive $\lambda$ (saddle point)
  - positive (and large) Laplacian: ionic bond
  - negative Laplacian: covalent bond
- bridge CP: (3,1)
- cage CP: (3,3) (minimum)

trajectories of constant $\nabla \rho$
originating at CPs in C$_2$H$_4$
“Atoms” are regions within a zero-flux surface $\nabla \rho \cdot \vec{n} = 0$

$\rho$ of $C_2H_4$ with zero-flux lines defining atomic basins
example of BN/Ni with “difference” to free atoms,
workfunction shift
Bader analysis of some inorganic compounds:

<table>
<thead>
<tr>
<th></th>
<th>( \rho(\text{e/A}^3) )</th>
<th>( \Delta \rho(\text{e/A}^5) )</th>
<th>Q (e)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Cl(_2)</td>
<td>1.12</td>
<td>-6.1</td>
<td>-</td>
</tr>
<tr>
<td>I(_2)</td>
<td>0.48</td>
<td>-0.9</td>
<td>-</td>
</tr>
<tr>
<td>TiC</td>
<td>0.51</td>
<td>1.8</td>
<td>1.7</td>
</tr>
<tr>
<td>TiN</td>
<td>0.47</td>
<td>3.9</td>
<td>1.7</td>
</tr>
<tr>
<td>TiO</td>
<td>0.43</td>
<td>5.8</td>
<td>1.5</td>
</tr>
<tr>
<td>KCl</td>
<td>0.08</td>
<td>1.2</td>
<td>0.6</td>
</tr>
</tbody>
</table>

Cl\(_2\) more covalent than I\(_2\)
more ionic, but less charge?
less ionic than TiC?
You must have a “good” scf-density (case.clmsum)

- no core leakage, LMs up to L=8-10 in case.in2

SURF
1 atom in center of surface (including MULT)
20 0.0 1.570796327 theta, 20 points, from zero to pi/2
20 0.0 0.785398163 phi, from 0 to pi/4 (depends on symmetry!!)
0.07 1.0 4 step along gradient line, rmin (has reached an atom)
1.65 0.1 initial R for search, step (a.u)
3 3 3 nshell
IRHO "INTEGRATE" rho
WEIT WEIT (surface weights are available in case.surf)
30 30 radial points outside min(RMIN,RMT)
END

---------------------CRIT
1 atom around you search for critical points
ALL two, three, four, all (dimers, trimers, ....all=2+3)
3 3 3 nshell
END

extractaim_lapw: ➔ critical_points_ang (converted units)
:PC x, y, z, \lambda_1, \lambda_2, \lambda_3, ch, laplacian, rho
Properties with WIEN2k - II

- **Total energy and forces**
  - optimization of internal coordinates, (MD, BROYDEN)
  - cell parameter only via $E_{\text{tot}}$ (no stress tensor)
  - elastic constants for cubic, hexagonal, and tetragonal cells
  - Phonons via supercells
    - interface to PHONON (K.Parlinski) – bands, DOS, thermodynamics, neutrons
    - interface to PHONOPY (A. Togo)
      - http://www.wien2k.at/reg_user/unsupported

- **Spectroscopy**
  - core level shifts
  - X-ray emission, absorption, electron-energy-loss (with core holes)
    - core-valence/conduction bands including matrix elements and angular dep.
  - optical properties (dielectric function in RPA approximation, JDOS including momentum matrix elements and Kramers-Kronig)

- **fermi surface:** 2D, 3D (using XcrysDen)
Fermi surfaces

- `xcrysden --wien_fermisurface tin.struct`
  - choose a good k-mesh (eg. 10000 points); (DON'T CHANGE to UNIT 5 !!!)
  - plot the FS for all bands which cross $E_F$ and compare to band structure

- for 2D plots there is also a WIEN2k-tool „fsgen“ (see UG)
- SKEAF ([www.wien2k.at/reg_users/unsupported](http://www.wien2k.at/reg_users/unsupported)): quantum oscillations
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- fermi surface: 2D, 3D *(using XcrysDen)*
Cohesive energy

\[ E_{\text{cohes.}}^{A_xB_y} = E_{\text{crystal}} - xE_{\text{atom}}^A - yE_{\text{atom}}^B \]

- **E\text{crystal}**: scalar-relativistic valence (or approx. SO)
- **E\text{atom}**: LSTART: fully-relativistic\(\rightarrow\) inconsistent description
  - \(\rightarrow\) for heavier elements (2\text{nd} row):
    - supercell with one atom in a \(\sim\)30 bohr FCC box
      (identical RMT, RKmax, 1 k-point, spin-polarized)
Structural optimizations:

- Lattice parameters, volume, c/a ratio only via total energies:
  - `x optimize`: creates a series of "struct" files + script "optimize.job"
    - select volume or c/a, ...
    - select number of cases and desired changes in volume (in % of $V_0$)
  - edit `optimize.job`
    - adapt to your need: change / uncomment various lines, eg.:
      - select different convergence parameters, parallelization, more iterations (-i 40)
      - different "save_lapw" (with more specific names)
      - replace "run_lapw" by "runsp_lapw" or `min_lapw -l -j "run_lapw -l -fc 1"`
  - execute `optimize.job`
  - plot (analyse) the results

- combinations of volume and c/a are possible: 2Doptimize
  - "x optimize" always uses `case_initial.struct` (if present)
  - do a "volume" optimization to create `case_vol_xx.struct` files
  - copy the respective `case_vol_xx.struct` file to `case_initial.struct`
  - `x optimize` with "c/a" for this particular volume and proceed as above.
Symmetry:

- WIEN „preserves“ symmetry:
  - \(c/a\) optimization of „cubic“ TiC:
    - change \(c\) lattice parameter in TiC.struct (tetragonal distortion, \#sym.op=0)
    - init_lapw
    - change \(c\) back to cubic
    - \(x\) optimize …
  - „Jahn-Teller“ distortion:
    - when you start with a perfect octahedra, you will never get any distortion
    - ➔ start with slightly distorted positions
Supercells

$2 \times 2 \times 2 = 8 \text{ atoms}$

(0,0,0)  P $\Rightarrow$ 8 atoms (0,0,0) (.5,0,0) (.5,.5,0) (.5,.5,.5) (0,.5,0) (.5,0,.5) (0,0,.5) (0,.5,.5) (0,0,.5) (0,.5,.5)

B $\Rightarrow$ 4 atoms yes yes no no

F $\Rightarrow$ 2 atoms yes no no yes

4x4x4 supercells:  P (64), B (32), F (16) atoms

$\sqrt{2} \times \sqrt{2}$ supercells (1 $\Rightarrow$ 2 atoms)
Supercells

- Program „supercell“:
  - start with „small“ struct file
  - specify number of repetitions in x,y,z (only integers, e.g. 2x2x1)
  - specify P, B or F lattice
  - add „vacuum“ for surface slabs (only (001) indexed surfaces)
  - shift all atoms in cell

- You must break symmetry!!!
  - replace (impurities, vacancies) or
  - displace (phonons) or
  - label at least 1 atom (core-holes, specific magnetic order; change “Fe” to “Fe1”; this tells the symmetry-programs that Fe1 is NOT a Fe atom!!)

- At present „supercell“ works only along unit-cell axes!!!
Structureditor (by R. Laskowski)

- requires octave (matlab) and xcrysden (visualization)
- allows complex operations on struct-files

```octave
s=loadstruct("GaN.struct")
# make an orthorhombic supercell and visualize it
a=[1 0 0; 1 1 0; 0 0 2]
sout=makesupercell(s,a);
showstruct(sout);
# save it as test.struct
savestruct(sout,"test.struct");
# get help on all commands
helpstruct
```
Surfaces

- 2D-slabs with finite number of layers with „vacuum“ in 3rd dimension

**bcc (001) 7 layers:**

- (0 0 6z) (.5 .5 +/-3z) with lattice parameters:
- (.5 .5 5z) (0 0 +/-2z) a, a, c=(3a+15-20bohr vacuum)
- (0 0 4z) shift to (.5 .5 +/-z)
- (.5 .5 3z) (0 0 0) z= a/2c
- (0 0 2z) inversion
- (.5 .5 z)
- (0 0 0)

**bcc (110):**

Orthorhombic CXY-lattice: a, $\sqrt{2}a$, c

- (0 0 0) $z=\frac{a}{\sqrt{2}a}$ c
- (0 .5 +/-z)
- (0 0 +/-2z)

$$\sqrt{2}a$$
Work function

potential

supercell

Surface

Vacuum

Work function

WF = ZER0 - FER (check convergence with vacuum)
Total energies and atomic forces

(Yu et al.; Kohler et al.)

■ **Total Energy:**
  - *Electrostatic energy*
  - *Kinetic energy*
  - *XC-energy*

**Force on atom \( \alpha \):**

\[
U[\rho] = \frac{1}{2} \int d^3\vec{r} \, \rho(\vec{r}) V_{es}(\vec{r}) + \frac{1}{2} \sum_{\alpha} Z_{\alpha} V_{es}^{\alpha}(\vec{r})
\]

\[
T[\rho] = \sum_i n_i \varepsilon_i - \int d^3\vec{r} \, \rho(\vec{r}) V_{eff}(\vec{r})
\]

\[
E_{xc}[\rho] = \int d^3\vec{r} \, \rho(\vec{r}) \varepsilon_{xc}(\vec{r})
\]

\[
\vec{F}^{\alpha} = -\frac{dE_{tot}}{d\vec{R}_\alpha} = F_{HF}^{\alpha} + F_{core}^{\alpha} + F_{val}^{\alpha}
\]

■ **Hellmann-Feynman-force**
  - *Core*
  - *Valence*

■ **Pulay corrections**
  - *Core*
  - *Valence*

■ expensive, contains a summation of matrix elements over all occupied states

\[
F_{HF}^{\alpha} = Z_{\alpha} \sum_{m=-1}^{1} \lim_{r_{\alpha} \rightarrow 0} \frac{V_{es}^{m}(r_{\alpha})}{r_{\alpha}} \nabla_{\alpha} [r_{\alpha} Y_{1m}(\vec{r})]
\]

\[
F_{core}^{\alpha} = -\int \rho_{core}(r) \nabla^{\alpha} V_{eff}(r) \, d\vec{r}
\]

\[
F_{val}^{\alpha} = \int_{\alpha} V_{eff}(r) \nabla^{\alpha} \rho_{val}(r) \, d\vec{r} + \sum_{K,K'} \sum_{k,i} c^{*}_{i}(K') c_{i}(K) \times
\]

\[
\left[ (K^2 - \varepsilon_i) \int \phi^{*}_{K'}(r) \phi_{K}(r) \, dS_{\alpha} - i(K - K') \langle \phi_{K'} | H - \varepsilon_i | \phi_{K} \rangle_{\alpha} \right]
\]
Optimization of internal parameters using “forces”

- Forces only for “free” structural parameters:
  - NaCl: (0,0,0), (0.5,0.5,0.5): all positions fixed by symmetry
  - TiO$_2$: Ti (0,0,0), O (u,u,0): one free parameter (u,x,y,z)

- Forces are only calculated when using “-fc”:
  - run_lapw -fc 1.0 (mRy/bohr)
    - grep :fgl002 case.scf
      - 200. partial
      - -130. partial
      - 140. partial
      - 135 partial only $F_{HF} + F_{core}$
      - 120 partial
      - 122 partial forces converging
      - 121 partial $\rightarrow$ changes “TOT” to “FOR” in case.in2
      - -12.3 total $F_{HF} + F_{core} + F_{val}$, only this last number is correct

- Forces are useful for
  - structural optimization (of internal parameters)
  - phonons
Structure optimization (atomic positions)

Traditional way:

- Inner loop: obtain fixed-point for given atom positions
- Outer loop: optimize atomic positions
Current algorithms

- Calculate SCF mapping, time $T_0$
- Broyden expansion for fixed-point problem, self-consistent density, $N_{SCF}$ iterations
- BFGS is most common for optimizing the atomic positions (Energy), $N_{BFGS}$
- Time scales as $N_{SCF} \times N_{BFGS} \times T_0$

Each step is a **full** scf calculation producing **accurate** forces

Structural optimization of internal parameters using “PORT”

- /home/pblaha/tio2> min_lapw [-p -it -sp] [-j “run -fc 1 -p -it”] [-NI]
  - performs scf-cycle for fixed positions
  - get forces and move atoms along forces (building an approximate Hessian) and writing a new case.struct file
  - extrapolate density (case.clmsum)
  - perform next scf cycle and loop until forces are below „tolf“
- CONTROL FILES:
  - .minstop stop after next structure change
- tio2.inM (generated automatically by “pairhess” at first call of min_lapw)
  - PORT 2.0 #(NEW1, NOSE, MOLD, tolf (a4,f5.2))
  - 0.0 1.0 1.0 1.0 # Atom1 (0 will constrain a coordinate)
  - 1.0 1.0 1.0 1.0 # Atom2 (NEW1: 1,2,3:delta_i, 4:eta (1=MOLD, damping))
- monitor minimization in file case.scf_mini
  - contains last iteration of each geometry step
  - each step N is saved as case_N.scf (overwritten with next min_lapw !)
    - grep :ENE case.scf_mini
    - grep :FGLxxx case.scf_mini (:POSxxx)
Optimization of atomic positions (E-minimization via forces)

- damped Newton mechanics scheme (NEW1: with variable step)
- quite efficient quasi-Newton (PORT) scheme
  - minimizes E (using forces as gradients and construct approx. Hessian)
  - If minimizations gets stuck or oscillates: (because E and \( F_i \) are inconsistent):
    - touch .minstop; min –nohess (or rm case.tmpM .min_hess)
    - improve scf-convergence (-ec), Rkmax, k-mesh, ...
    - change to NEW1 scheme

W impurity in Bi (2x2x2 supercell: Bi\(_{15}\)W)
Alternative method: **Fused Loop**

- Treat the **density** and **atomic positions all** at the same time.

- No restrictions to “special” cases, general algorithm has to work for insulators, metals, semiconductors, surfaces, defects, hybrids….

- Few to no user adjustable parameters

---

each step is a **single** scf cycle producing only **approximate** forces

Broyden Fixed-Point Methods

- Solve $(\rho(r,x) - F(\rho(r,x)), G) = 0$
- $s_k = (\rho, x)_{k+1} - (\rho, x)_k; y_k = (F(\rho, x), G)_{k+1} - (F(\rho, x), G)_k$
- Broyden’s “Good Method”

$$B_{k+1} = B_k + \frac{(y_k - B_k s_k) s_k^T}{s_k^T s_k} \quad H_{k+1} = H_k + \frac{(s_k - H_k y_k) s_k^T}{s_k^T y_k}$$

- Broyden’s “Bad Method”

$$H_{k+1} = H_k + \frac{(s_k - H_k y_k) y_k^T}{y_k^T y_k}$$

- Generalizable to multisecant method (better,


Comparison of the 2 methods

Larger Problems:
- 52 atoms, MgO (111) + H₂O
- 108 atoms AlFe


Lyudmila V. Dobysheva (2011)

- edit case.inm and set „MSR1a“
- run_lapw -fc 1.0 -cc 0.001 -ec 0.0001 [-it -noHinv -p ]

This runs ONE big scf-calculations optimizing the density and the positions (forces towards zero) simultaneously (may need hundreds of iterations).

- Monitor: :ENE and :FR (av. and max forces, movements)

- it continues until all :FR quantities are below „tolf“ (case.inM) and switches then automatically to MSR1 for a final charge optimization (with fixed positions).

- quite efficient, recommended method, still under development by L. Marks (Northwestern Univ).
Calculations of Phonons: The Direct Method

WIEN2k + Phonon

Copyright by K. Parlinski

http://wolf.ifj.edu.pl/phonon/

alternatively use A. Togo's PHONOPY code + Wien2k-interface
(see www.wien2k.at unsupported)
THEORY OF DIRECT METHOD

System energy $E$ (at $T = 0$) as a function of atomic positions $\mathbf{R}(n, \mu)$ is

$$E(\mathbf{R}(n, \mu), \ldots \mathbf{R}(m, \nu), \ldots) = E_0 + \frac{1}{2} \sum_{n,\mu,m,\nu} \Phi(n, \mu, m, \nu) U(n, \mu) U(m, \nu)$$

where the force constant matrix are

$$\Phi_{i,j}(n, \mu, m, \nu) = \left. \frac{\partial^2 E}{\partial \mathbf{R}_i(n, \mu) \partial \mathbf{R}_j(m, \nu)} \right|_0$$

is defined at $\frac{\partial E}{\partial \mathbf{R}_i(n, \mu)} |_0 = 0$.

The dynamical matrix is defined as

$$D(k; \mu, \nu) = \frac{1}{\sqrt{M_{\mu} M_{\nu}}} \sum_m \Phi(0, \mu; m, \nu) \exp \left\{ -2\pi i k \cdot [\mathbf{R}(0, \mu) - \mathbf{R}(m, \nu)] \right\}$$

$m$ runs over all atoms. Diagonalization of the dynamical matrix

$$\omega^2(k, j) \mathbf{e}(k, j) = D(k) \mathbf{e}(k, j)$$

gives phonon frequencies $\omega^2(k, j)$ and polarization vectors $\mathbf{e}(k, j)$.

Any atomic displacement $U(m, \nu)$ generates forces

$$F(n, \mu) = -\frac{\partial E}{\partial \mathbf{R}(n, \mu)}$$

on all other atoms. Hence

$$F_j(n, \mu) = -\sum_{m,\nu,j} \Phi_{i,j}(n, \mu, m, \nu) U_j(m, \nu)$$

Master equation of direct method.
CUMMULANT FORCE CONSTANTS

Displace an atom by \( U(m, \nu) \)

\[
F_i(n, \mu) = - \sum_L \Phi_{i,j}(n, \mu, m + L, \nu) U_j(m, \nu)
\]

\( L = (L_a, L_b, L_c) \) are the indices of supercell lattice constants.
or

\[
F_i(n, \mu) = -\Phi^{(\Sigma)}_{i,j}(n, \mu, m, \nu) U_j(m, \nu)
\]

where the **cummulant force constant** is

\[
\Phi^{(\Sigma)}_{i,j}(n, \mu, m, \nu) = \sum_L \Phi_{i,j}(n, \mu, m + L, \nu)
\]

\( L \) runs over all supercell images.
Supercell dynamical matrix. Exact wave vectors.

Conventional dynamical matrix:

\[
D(k; \mu, \nu) = \frac{1}{\sqrt{M_\mu M_\nu}} \sum_m \Phi(0, \mu; m, \nu) \exp\{-2\pi i k \cdot [R(0, \mu) - R(m, \nu)]\}
\]

Supercell dynamical matrix:

\[
D^{(SC)}(k; \mu, \nu) = \frac{1}{\sqrt{M_\mu M_\nu}} \sum_{m \in SC} \Phi^{(SC)}(0, \mu; m, \nu) \exp\{-2\pi i k \cdot [R(0, \mu) - R(m, \nu)]\}
\]

These two matrices are equal if

\[
D^{(SC)}(k; \mu, \nu) = D(k; \mu, \nu)
\]

- **interaction range** is confined to interior of supercell (supercell is big enough)
- wave vector is **commensurate with the supercell** and fulfils the condition (independent of interaction range):

\[
\exp\{-2\pi i k_s \cdot L\} = 1
\]

At wave vectors \(k_s\) the phonon frequencies are “exact”, provided the supercell contains the complete list of neighbors.

Wave vectors \(k_s\) are commensurate with the supercell size.
Exact wave vectors

1x1x1

Exact: $\Gamma$

2x2x2

Exact: $\Gamma, X, M, R$

3x3x3

Exact: $\Gamma$

$\text{Exact: } \Gamma$
Phonon dispersions + density of states

$\omega$

Total + Germanium

$\omega$

Total + Oxygen

GeO$_2$  P4$_2$/mnm
Thermodynamic functions of phonon vibrations

Internal energy:
\[ E = \frac{1}{2} r \int_0^\infty d\omega \, g(\omega) \left( \frac{\hbar \omega}{2k_B T} \right) \coth \left( \frac{\hbar \omega}{2k_B T} \right) \]

Free energy:
\[ F = r k_B T \int_0^\infty d\omega \, g(\omega) \ln \left[ 2 \sinh \left( \frac{\hbar \omega}{2k_B T} \right) \right] \]

Entropy:
\[ S = r k_B \int_0^\infty d\omega \, g(\omega) \left\{ \left( \frac{\hbar \omega}{2k_B T} \right) \coth \left( \frac{\hbar \omega}{2k_B T} \right) - 1 \right\} - \ln \left[ 1 - \exp \left( - \frac{\hbar \omega}{k_B T} \right) \right] \]

Heat capacity \( C_V \):
\[ C = r k_B \int_0^\infty d\omega \, g(\omega) \left( \frac{\hbar \omega}{k_B T} \right)^2 \frac{\exp \left( \frac{\hbar \omega}{k_B T} \right)}{\left[ \exp \left( \frac{\hbar \omega}{k_B T} \right) - 1 \right]^2} \]

Thermal displacements:
\[ B_{ij}(\mu) = \langle U_i(\mu) U_j(\mu) \rangle \]
\[ B_{ii}(\mu) = \frac{\hbar r}{2M_\mu} \int_0^\infty d\omega \, g_{ii,\mu}(\omega) \frac{1}{\omega} \coth \left( \frac{\hbar \omega}{2k_B T} \right) \]
PHONON

- by K.Parlinski (Crakow)
- Linux or MS-windows
- uses a „direct“ method to calculate Force-constants with the help of an ab initio program
- with these Force-constants phonons at arbitrary k-points can be obtained

- Define your spacegroup
- Define all atoms

http://wolf.ifj.edu.pl/phonon/
Phonons:

- select symmetry adapted atomic displacements (4 displacements in cubic perovskites)

- select a supercell: (eg. 2x2x2 atom P-type cell)

- calculate all forces for these displacements with high accuracy (WIEN2k)

  - force constants between all atoms in the supercell
  - dynamical matrix for arbitrary q-vectors
  - phonon-dispersion ("bandstructure") using PHONON (K.Parlinski)
- Define an interaction range (supercell)
  - create displacement file
  - transfer case.d45 to Unix
- Calculate forces for all required displacements
  - init_phonon_lapw
    - for each displacement a case_XX.struct file is generated in an extra directory
    - runs nn and lets you define RMT values like:
      - 1.85 1-16

- init_lapw: either without symmetry (and then copies this setup to all case_XX)
  or with symmetry (must run init_lapw for all case_XX) (Do NOT use SGROUP)
- run_phonon: run_lapw -fc 0.1 -i 40 for each case_XX
PHONON-III

- **analyze_phonon_lapw**
  - reads the *forces* of the scf runs
  - generates „Hellman-Feynman“ file *case.dat* and a „symmetrized HF-file case.dsy“ (when you have *displacements in both directions*)
    - check quality of forces:
    - sum $F_x$ should be small (0)
    - $\text{abs}(F_x)$ should be similar for +/- displacements
- transfer *case.dat* (dsy) to Windows
- Import HF files to PHONON
- Calculate force constants
- Calculate phonons, analyze phonons eigenmodes, thermodynamic functions
Applications:

- phonon frequencies (compare with IR, raman, neutrons)
- identify dynamically unstable structures, describe phase transitions, find more stable (low T) phases.

Pyrochlore structure of $Y_2Nb_2O_7$: strong phonon instabilities $\rightarrow$ phase transition
from E. Assmann, TU Vienna

needs **wannier90** !

*Wannier90: A Tool for Obtaining Maximally-Localised Wannier Functions*

A. A. Mostofi, J. R. Yates, Y.-S. Lee, I. Souza, D. Vanderbilt and N. Marzari

Bloch functions - Wannier functions

Bloch functions

\[ \psi_{nk}(r) = u_{nk}(r)e^{ik \cdot r} \]

Indexed by the wave vector

Wannier functions

\[ |R_n\rangle = \frac{V}{(2\pi)^3} \int_{BZ} dke^{-i(k \cdot R) \psi_{nk}} \]

Indexed by the lattice vector in real space

\[ \psi_{k_0}(x) \]

\[ \psi_{k_1}(x) \]

\[ \psi_{k_2}(x) \]

Both sets: complete and orthonormal

Maximally Localized Wannier Functions

Bloch functions (more precisely):

$$\psi_{nk}(\mathbf{r}) = w_{nk}(\mathbf{r}) e^{i\mathbf{k} \cdot \mathbf{r}} e^{i\phi_{nk}}$$

gauge freedom

(does not change the physical description of the system)

$$\langle r^2 \rangle_n - \bar{r}^2_n$$

$$|w_{nR}\rangle = \frac{V}{(2\pi)^3} \int_{BZ} d\mathbf{k} \sum_m e^{i\mathbf{k} \cdot \mathbf{R}} U_{mn}(\mathbf{k}) |\psi_{m\mathbf{k}}\rangle$$

$$\Omega = \sum_n [\langle 0 n | r^2 | 0 n \rangle - \langle 0 n | r | 0 n \rangle^2] = \sum_n [\langle r^2 \rangle_n - \bar{r}^2_n]$$

choose $U(\mathbf{k})$ to minimize $\Omega$
maximally localized wannier functions

- choose $U(k)$ to minimize spread $\rightsquigarrow$ MLWF
- total spread $\Omega = \Omega_1 + \tilde{\Omega}$ can be split into gauge-invariant part and rest
  $\rightsquigarrow$ minimize $\tilde{\Omega}$

- wannier90 computes $U(k)$ in this way

  [Marzari et al., Rev. Mod. Phys. 84, 1419 (2012)]
  [http://wannier.org]

- wien2wannier provides interface to Wien2k

  [Kuneš, Wissgott et al., Comp. Phys. Commun. 181, 1888]
Disentanglement

- other bands may cross target manifold

\[ \rightarrow \text{must select bands to Wannierize} \]
  - \[ V(k) \left( j(k) \times N \right) \]

- selection determines \( \Omega_1 \)

\[ \rightarrow \text{minimize also } \Omega_1 \]

from Marzari et al.

**fcc-Cu, 5 d-like WF, 2 interstitial s-like WF**

\text{num\_bands} = 12

\text{num\_wann} = 7
MLWF: applications

- analysis of chemical bonding
- electric polarization and orbital magnetization
  \[ \Rightarrow \text{Oleg Rubel’s talk} \]
- Wannier interpolation \( \mathcal{K} \rightarrow \mathcal{G} \)
  \[ H(k)|_{\mathcal{K}} \xrightarrow{\mathcal{F}} H(R)|_{\mathcal{K}^{-1}} \xrightarrow{\mathcal{F}^{-1}} H(k)|_{\mathcal{G}} \]
- Wannier functions as basis functions
  - tight-binding model \( H(k) = U^+(k) \varepsilon(k) U(k) \)
  \[ \Rightarrow \text{realistic dynamical mean-field theory (DMFT)} \]
Wannier functions as tight-binding basis

\[$\text{less GaAs-WANN_hr.dat}$\]

\[
\begin{array}{cccc}
0 & 0 & 0 & 1 \\
0 & 0 & 0 & 2 \\
0 & 0 & 0 & 3 \\
0 & 0 & 0 & 4 \\
0 & 0 & 0 & 5 \\
0 & 0 & 0 & 6 \\
0 & 0 & 0 & 7 \\
0 & 0 & 0 & 8 \\
\end{array}
\]

\[\langle s_1 | s_1 \rangle\]

\[\langle s_1 | H | s_1 \rangle = E_{s_1}\]

\[0.000000 \text{ Im part} = 0\]

\[\langle s_2 | H | s_1 \rangle = V_{ss\sigma}\]

\[\langle p_2 | H | s_1 \rangle = V_{sp}\]

Home unit cell

Neighbour unit cell

WF are well localized \(\Rightarrow\) nearest-neighbour suffice
workflow using wien2wannier

- Regular SCF calculation
- Band structure plot

- Initialize wien2wannier (init_w2w):
  - select bands, init. projections, # of WF (case.inwf file)
  - projected band structure “bands_plot_project” (case.win file)
  - additional options related to entanglement (case.win file)
- Compute overlap matrix element $M_{mn}, A_{mn}$ (x w2w)

- Perform Wannierization (x wannier90):
  - position of Wannier centers and spreads (case.wout file)
  - Wannier hamiltonian (case_hr.dat file)

- Initialize plotting, select plotting range, r-mesh (write_inwpplot)
- Evaluate WF on the r-mesh selected (x wplot)
- Convert the output of wplot into xcrysden format for plotting (wplot2xsf)

Relation to Polarization

Bond-centered WF

Si

GaAs

symmetric

non-symmetric

\[ P = \frac{e}{V} \left( \sum_{\tau} Z_{\tau} r_{\tau} - \sum_{n} r_{n} \right) \]

Ionic part

Electronic part

+ \( Z_{As} \)

- \( Z_{Ga} \)

Wannier center \( q_e \)

from O. Rubel (Lakehead University)
needs wien2wannier, Python 2.7.x, numpy
Materials properties related to Polarization

Piezoelectricity

Dielectric screening

Effective charge

Pyroelectricity

"Fire" (i.e., Electricity)
What is Polarization?

We will now assume that in each atom there are charges $q$ separated by a distance $\delta$, so that $q\delta$ is the dipole moment per atom. (We use $\delta$ because we are already using $d$ for the plate separation.) If there are $N$ atoms per unit volume, there will be a dipole moment per unit volume equal to $Nq\delta$. This dipole moment per unit volume will be represented by a vector, $P$. Needless to say, it is in the direction of the individual dipole moments, i.e., in the direction of the charge displacement.

Polarization for periodic solids is undefined.
Modern theory of polarization (Berry phase)

Pioneered by King-Smith, David Vanderbilt and Raffaele Resta

All measurable physical quantities are related to the change in polarization!

\[ \Delta P = P^{(0)} - P^{(1)} \]

- \[ \frac{\Delta P}{\Delta \text{strain}} \]
- \[ \frac{\Delta P}{\Delta E} \]
- \[ \frac{\Delta P}{\Delta T} \]

"Fire" (Pyro) Electricity
Polarization as a transient current

\[ \Delta P = P^{(1)} - P^{(0)} = \Omega^{-1} \int dt \int_{\text{cell}} dr \, j(r, t) \]

transient current density

\[ \mathbf{P} = \mathbf{P}_{\text{ion}} + \mathbf{P}_{\text{el}} \]

\[ \mathbf{P}_{\text{ion}} = \frac{e}{\Omega} \sum_s Z_s^{\text{ion}} \mathbf{r}_s \]

In Wien2k \( Z_s^{\text{ion}} \) is the core charge

\[ -\mathbf{P}_{\text{el}} = \Omega^{-1} \int dr \, \mathbf{r} \rho(r) = \Omega^{-1} \sum_{n, \text{occ., bands}} \left\langle \psi_n^* | \mathbf{r} | \psi_n \right\rangle \equiv \frac{2e\text{i}}{(2\pi)^3} \sum_n \left\langle \psi_n^* | \nabla | \psi_n \right\rangle \int_{\text{BZ}} dk \left\langle u_{nk}^* | \nabla_k | u_{nk} \right\rangle \]

King-Smith and David Vanderbilt, Phys. Rev. B 47, 1651 (1993)
Berry phase

\[ d\varphi_n = -i \langle u_{nk} | \nabla_k | u_{nk} \rangle \cdot dk = -i \ln \langle u_{nk} | u_{n(k+dk)} \rangle \]

\[ S_{mn}(k_j, k_{j+1}) = \langle u_{mk_j} | u_{nk_{j+1}} \rangle \quad \text{WIEN2WANNIER} \]

\[ \varphi(k_{||}) = 2 \text{Im} \left[ \ln \prod_{j=0}^{J-1} \det S_{MM}(k_j, k_{j+1}) \right] \]

\[ \varphi_{el,\alpha} = S_{\perp}^{-1} \int_{S_{\perp}} dS_{\perp} \varphi(k_{||}) \]

\[ P_\alpha = \frac{e(\varphi_{el,\alpha} + \varphi_{ion,\alpha})}{2\pi\Omega} R_\alpha \]
Polarization change for 2 cases:

\[ \Delta P = P^{(0)} - P^{(1)} \]

- structure file **must** preserve the symmetry
- begin with the lowest symmetry (\(\lambda_1\)) case
- copy case \(\lambda_1\) to case \(\lambda_0\)
- edit structure file for case \(\lambda_0\)
- do **not** initialize calculation (\textit{init\_lapw})
- update density (\textit{x\_dstart})
- run SCF cycle (\textit{run\_lapw})
- run Berry\textsc{PI}
Uncertainties

\[ P_\alpha = \frac{e(\varphi_{el,\alpha} + \varphi_{ion,\alpha})}{2\pi \Omega} R_\alpha \]

\[ \Delta P = P^{(0)} - P^{(1)} \pm \frac{e}{\Omega} R \]

- cannot determine large polarization difference
  \( \sim 1 \text{ C/m}^2 \)

Solution: \( \lambda_1 \Rightarrow \lambda_{1/2} \Rightarrow \lambda_1 \)
Spontaneous polarization in BaTiO$_3$

\[ P_S = P_{nc} - P_C \]
Born effective charge is related to polarization

$$Z^*_{s,\alpha\beta} = \frac{\Omega}{e} \frac{\partial P_{\alpha}}{\partial r_{s,\beta}}$$

Introduce small displacements
$$\pm \Delta r \ll a_0$$

P - Polarization
r - atom position
$$\Omega$$ - unit cell volume
e - elementary charge

Need to compute the polarization difference between 2 structures:
$$dP = P(+\Delta r) - P(-\Delta r)$$
Piezoelectricity

\[ \lambda_0 \]

\[ \lambda_1 \]

tetragonal equilibrium

\[ a_0 = b_0 \neq c_0 \]

tetragonal strained

\[ a_0 = b_0 \neq c_0(1+\varepsilon_3) \]